

## Analysis of the Tuggle Front End – Part II

We shall now consider the Tuggle tuner delivering power to a load. First, we must account for the parallel RF losses of the unloaded tuned circuits. Let  $R_p$  represent the losses of the tank circuit comprising  $L_1$  and  $C$ , and  $R_{TNK}$  those of  $L_2$  and  $C_2$  (please see Fig. 1).

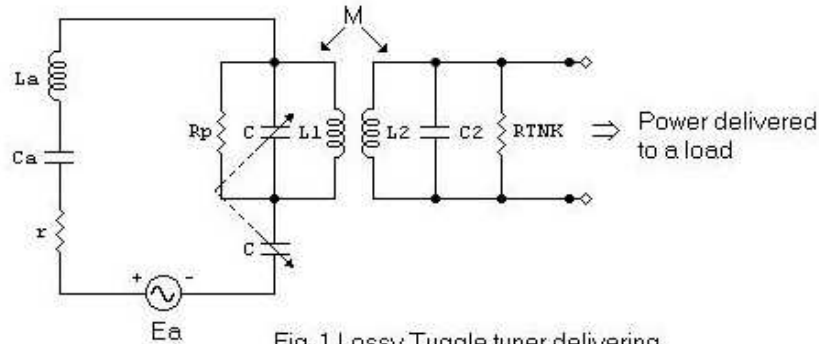


Fig. 1 Lossy Tuggle tuner delivering power to the secondary load

The load consists of a diode detector  $D$  in series with an audio load  $R_L$  (usually an audio transformer matching a pair of 2k ohms DC resistance magnetic headphones or low-impedance sound powered phones to the detector) and is coupled to the tuner via the magnetic coupling existing between  $L_1$  and  $L_2$ , being  $M$  the mutual inductance of the coils. The secondary is tuned to the same radian frequency as the primary. An schematic for the load can be seen in Fig. 2.

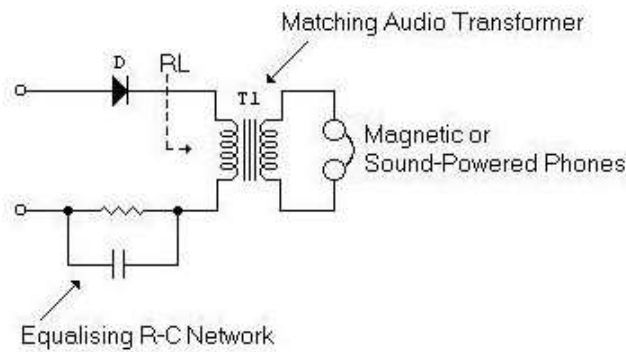


Fig. 2 Load coupled to the Tuggle tuner

Usually, it is assumed that optimum RF power transfer occurs when the antenna-ground system resonance resistance  $r$  is matched to the unloaded-secondary resonance parallel resistance  $R_{TNK}$ , with the diode detector's input resistance matched to this combination. Thus,

$$\frac{1}{R_{OPT}} = \frac{1}{R_{TNK}} + \frac{1}{R_{TNK}} + \frac{1}{\left(\frac{R_{TNK}}{2}\right)}$$

$$\text{or } R_{OPT} = \frac{R_{TNK}}{4}.$$

This is the overall parallel RF resistance of the secondary tank under matched conditions, and suggests that the unloaded Q of this tank circuit has been reduced to ¼ of its value.

In this case, the net parallel resistance to be coupled to the primary is:

$$R_2 = \frac{R_{TNK} \left( \frac{R_{TNK}}{2} \right)}{R_{TNK} + \left( \frac{R_{TNK}}{2} \right)}$$

$$= \frac{R_{TNK}}{3} \quad \dots(1)$$

We shall work on this later.

### Some circuit equivalents

In Fig.1, let's replace  $L_1$  and the coupled secondary circuit by the equivalent shown in Fig.3.a, which in turn can be replaced by the transformer circuit shown in Fig.3.b.

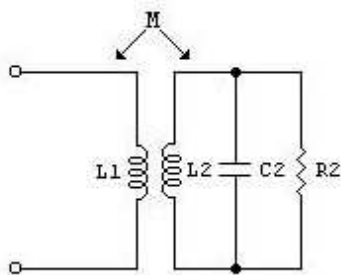
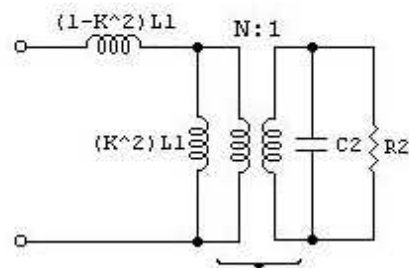


Fig.3.a Magnetically coupled circuits



Ideal Transformer  
 $N = K(L_1/L_2)^{0.5}$   
 $K = \text{Coupling Coefficient}$

Fig.3.b Equivalent transformer circuit

In the transformer circuit, the impedance coupled to the primary side consists of a capacitance  $C_2 / N^2$  in parallel with a resistance  $N^2R_2$ . Both are in parallel with the magnetizing inductance  $K^2L_1$  (please see Fig.4.a).

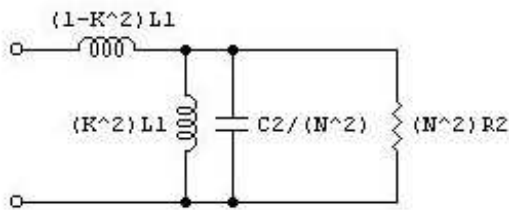


Fig.4.a Equivalent circuit as seen from the primary side

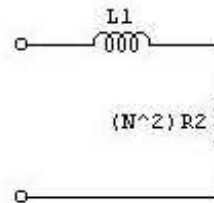


Fig.4.b Reduced equivalent at  $\omega = \omega_r$

$K^2L_1$  and  $C_2 / N^2$  resonate at a frequency

$$\omega_0 = \frac{1}{\sqrt{K^2L_1\left(\frac{C_2}{N^2}\right)}}$$

$$= \frac{1}{\sqrt{L_2C_2}} = \omega_r$$

as  $N^2 = K^2 \frac{L_1}{L_2}$ .

The equivalent circuit of Fig.4.a reduces to that of Fig.4.b, taking into account that for crystal set use, normally  $K \ll 1$ .

Up to this point, in the tuner side we have the equivalent circuit depicted in Fig.5.

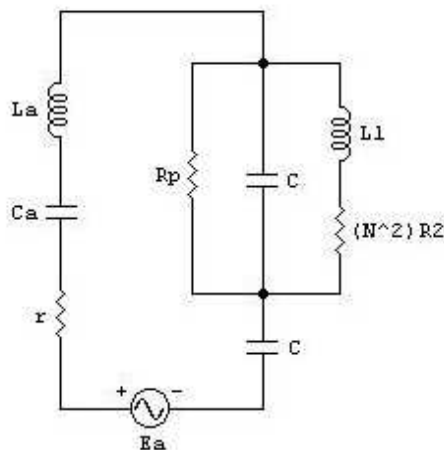


Fig.5 Equivalent circuit of the tuner with coupled secondary load

The series-coupled resistance  $N^2R_2$  can be transformed into a resistance  $R_{p1}$  in parallel with  $R_p$ . Using the known series-to-parallel “loss resistance” transformation we get:

$$R_{p1} = \frac{(\omega_r L_1)^2}{N^2 R_2} \quad \dots(2)$$

Let  $Q_2 = \frac{R_2}{\omega_r L_2}$ . Then:

$$R_{p1} = \frac{\omega_r L_1}{K^2 Q_2} \quad \dots(3)$$

Next, we compute the equivalent resistance  $R_T$  of the parallel combination of  $R_p$  and  $R_{p1}$ . It is given by:

$$R_T = \frac{R_p R_{p1}}{R_p + R_{p1}}$$

Substituting  $R_{p1}$  by its equivalent given by eq.(3):

$$\begin{aligned} R_T &= \frac{R_p \left( \frac{\omega_r L_1}{K^2 Q_2} \right)}{R_p + \frac{\omega_r L_1}{K^2 Q_2}} \\ &= \frac{R_p \omega_r L_1}{\omega_r L_1 + K^2 Q_2 R_p} \end{aligned}$$

Letting  $Q_1 = \frac{R_p}{\omega_r L_1}$  (unloaded Q of  $L_1$ -C tank) we obtain:

$$R_T = \frac{R_p}{1 + K^2 Q_1 Q_2} \quad \dots(4)$$

If  $R_T \gg \omega_r L_{eq} = \frac{\omega_r L_1}{1 - \omega_r^2 L_1 C}$  (please refer to part I of this study) we can redraw the equivalent circuit of the tuner at resonance as indicated by Fig.6.a, and, by virtue of the above inequality, the resonant frequency  $\omega_r$  will still be given by:

$$\omega_r^2 L_1 C \left( \frac{2Ca + C}{Ca + C} \right) = 1 \quad \dots(5)$$

Applying a parallel-to-series transformation, the equivalent circuit of Fig.6.b is obtained.

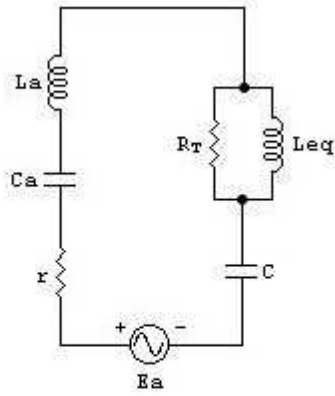


Fig.6.a Equivalent circuit of the tuner when  $R_T \gg \omega_r L_{eq}$

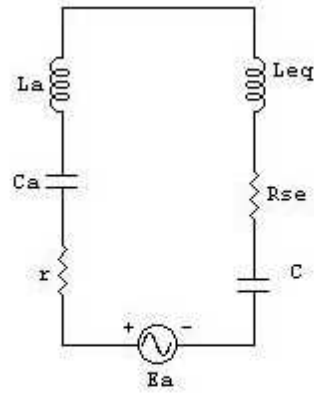


Fig.6.b Equivalent circuit after transformation

The series transformed resistance  $R_{se}$  is given by:

$$\begin{aligned}
 R_{se} &= \frac{(\omega_r L_{eq})^2}{R_T} \\
 &= (\omega_r L_{eq})^2 \left( \frac{1 + K^2 Q_1 Q_2}{R_p} \right) \\
 &= \frac{(\omega_r L_{eq})^2}{R_p} + (\omega_r L_{eq})^2 \frac{K^2 Q_1 Q_2}{R_p}
 \end{aligned}$$

Let  $R_s = \frac{(\omega_r L_{eq})^2}{R_p}$  and  $R_{s1} = (\omega_r L_{eq})^2 \frac{K^2 Q_1 Q_2}{R_p}$ .

$R_s$  is the series term due to  $R_p$  and  $R_{s1}$ , that coming from the coupled resistance  $N^2 R_2$ . Now, maximum RF power transfer to  $N^2 R_2$  occurs when:

$$r + R_s = R_{s1}$$

or when:

$$r + \frac{(\omega_r L_{eq})^2}{R_p} = (\omega_r L_{eq})^2 \frac{K^2 Q_1 Q_2}{R_p}$$

or

$$\frac{(\omega_r L_{eq})^2}{r} = \frac{R_p}{K^2 Q_1 Q_2 - 1} \quad \dots(6)$$

From part I of this study we know that:

$$\begin{aligned} L_{eq} &= \frac{L_1}{1 - \omega_r^2 L_1 C} \\ &= L_1 \left( 2 + \frac{C}{Ca} \right) \end{aligned} \quad \dots(7)$$

being

$$\omega_r^2 = \frac{1}{L_1 C} \left( \frac{Ca + C}{2Ca + C} \right)$$

Then:

$$\begin{aligned} (\omega_r L_{eq})^2 &= \frac{1}{L_1 C} \left( \frac{Ca + C}{2Ca + C} \right) L_1^2 \left( \frac{2Ca + C}{Ca} \right)^2 \\ &= \frac{L_1}{C} \left( 2 + \frac{C}{Ca} \right) \left( 1 + \frac{C}{Ca} \right) \end{aligned}$$

Eq.(6) is now written as:

$$\frac{L_1}{rC} \left( 2 + \frac{C}{Ca} \right) \left( 1 + \frac{C}{Ca} \right) = \frac{R_p}{K^2 Q_1 Q_2 - 1} \quad \dots(8)$$

Letting  $A = \frac{L_1}{rC} \left( 2 + \frac{C}{Ca} \right) \left( 1 + \frac{C}{Ca} \right)$ , eq.(8) takes the more compact form:

$$A = \frac{R_p}{K^2 Q_1 Q_2 - 1}$$

Solving for K we obtain:

$$K = \sqrt{\frac{R_p + A}{Q_1 Q_2 A}} \quad \dots(9)$$

which gives the value of the coupling coefficient for maximum RF power transfer to the secondary.

### **Power calculations**

The RF power delivered to the secondary load of Fig.1 will be at a maximum at resonance when eq.(6) is satisfied, this is, when  $r + R_s = R_{s1}$ . The maximum available power is then:

$$\begin{aligned}
P_{MAX} &= \frac{\left(\frac{E_a}{2}\right)^2}{2R_{s1}} \\
&= \frac{E_a^2}{8R_{s1}} \quad \dots(10)
\end{aligned}$$

where  $E_a$  is the peak value of the voltage induced in the antenna. The power delivered to the secondary load  $R_2$  is the same as that dissipated by the coupled resistance  $N^2R_2$ . To compute this power we need the voltage across  $L_1$  at resonance. This is the same as the voltage across  $L_{eq}$  in Fig.6.a. Then:

$$E_{Leq} = \frac{E_a}{2R_{s1}} j\omega_r L_{eq} \quad \dots(11)$$

From Fig.5 we obtain for the voltage across  $N^2R_2$ :

$$E_2' = \frac{E_{Leq}}{j\omega_r L_1 + N^2R_2} N^2R_2$$

In crystal sets,  $N^2R_2 \ll \omega_r L_1$ , due to the loose coupling between  $L_1$  and  $L_2$ . Then:

$$\begin{aligned}
E_2' &= \frac{E_{Leq}}{j\omega_r L_1} N^2R_2 \\
&= \frac{E_{Leq}}{j\omega_r L_1} \left[ \frac{(\omega_r L_1)^2}{R_{p1}} \right] \\
&= \frac{E_{Leq}}{jR_{p1}} \omega_r L_1
\end{aligned}$$

Bearing in mind eq.(3):

$$E_2' = \frac{E_{Leq}}{j} K^2 Q_2$$

Substituting the value of  $E_{Leq}$  given by eq.(11) into the above expression we obtain:

$$E_2' = \frac{E_a}{2R_{s1}} \omega_r L_{eq} K^2 Q_2$$

We can recall that:

$$R_{s1} = (\omega_r L_{eq})^2 \frac{K^2 Q_1 Q_2}{R_p} \quad \dots(12)$$

Then:

$$\begin{aligned} \frac{\omega_r L_{eq} K^2 Q_2}{R_{s1}} &= \frac{R_p}{\omega_r L_{eq} Q_1} \\ &= \frac{L_1}{L_{eq}} \\ &= \frac{Ca}{2Ca + C} \end{aligned}$$

Then, we obtain:

$$E_2' = \frac{E_a}{2} \left( \frac{Ca}{2Ca + C} \right)$$

The power dissipated by  $N^2 R_2$  is:

$$\begin{aligned} P_2 &= \frac{(E_2')^2}{2N^2 R_2} \\ &= \frac{\left( \frac{E_a}{2} \right)^2 \left( \frac{Ca}{2Ca + C} \right)^2}{2N^2 R_2} \quad \dots(13) \end{aligned}$$

Eq.(12) can be written as follows:

$$\begin{aligned} R_{s1} &= (\omega_r L_{eq})^2 \left( \frac{N^2 L_2}{L_1} \right) \left( \frac{Q_1 Q_2}{R_p} \right) \\ &= (\omega_r L_{eq})^2 \left( \frac{N^2 L_2}{L_1} \right) \left( \frac{1}{\omega_r L_1} \right) \left( \frac{R_2}{\omega_r L_2} \right) \\ &= \left( \frac{L_{eq}}{L_1} \right)^2 N^2 R_2 \end{aligned}$$

$$= \left( \frac{2Ca + C}{Ca} \right)^2 N^2 R_2$$

or:

$$N^2 R_2 = \left( \frac{Ca}{2Ca + C} \right)^2 R_{s1}$$

Substituting this equivalence into eq.(13):

$$P_2 = \frac{E_a^2}{8R_{s1}} = P_{MAX}$$

according to eq.(10).

$P_{MAX}$  is then dissipated by  $N^2 R_2$  and by consequence, this power is delivered to the secondary load.

### Some experimental results

Two coils,  $L_1$  and  $L_2$ , were wound on 4.5" diameter styrene forms using 660/46 Litz wire.  $L_1$  measured 152 uH and  $L_2$ , 222 uH. A two-gang 475 pF variable capacitor with bakelite insulation was used to tune  $L_1$ .  $L_2$  was tuned with a 480 pF variable capacitor with ceramic insulators.

Unloaded  $Q_s$  for each of the tuned circuits were measured at three frequencies. Accordingly, the corresponding RF losses were calculated. Data is tabulated below.

f	$Q_1$	$Q_{2UL}$	$Q_2$	$R_{pcalc}$	$R_{TNKcalc}$	C	$C_2$
530kHz	700	810	270	354.32 kohms	598.816 kohms	450 pF	406 pF
1 MHz	585	630	210	558.70 kohms	878.766 kohms	100 pF	114 pF
1.7 MHz	300	318	106	487.072 kohms	754.065 kohms	31 pF	39.5 pF

C : two-gang 475 pF variable capacitor with bakelite insulation

$C_2$  : 480 pF variable capacitor with ceramic insulation

$Q_{2UL}$  : unloaded Q of  $L_2$ - $C_2$  combination

$Q_1, Q_2$  : defined in the text

$R_p, R_{TNK}$ : defined in the text

Using the tabulated data, values for the optimum coupling coefficient K will be calculated for a working crystal set.

**f = 530 kHz**

$L_1 = 152$  uH

$Ca = 200$  pF (assumed)

$r = 30$  ohms (assumed)

$A = 1.555 \times 10^5$  ohms

$Q_1 Q_2 A = 2.938 \times 10^{10}$  ohms

$K = 4.165 \times 10^{-3}$

Check:  $\frac{R_p}{1 + K^2 Q_1 Q_2} = 82.8 \text{ kohms} \gg \omega_r L_{eq} = 2.214 \text{ kohms}$

**f = 1 MHz**

$L_1 = 152$  uH

$Ca = 200$  pF (assumed)

$r = 30$  ohms (assumed)

$A = 1.9 \times 10^5$  ohms

$Q_1 Q_2 A = 2.334 \times 10^{10}$  ohms

$K = 5.663 \times 10^{-3}$

$$\text{Check: } \frac{R_p}{1 + K^2 Q_1 Q_2} = 113.1 \text{ kohms} \gg \omega_r L_{eq} = 2.387 \text{ kohms}$$

**f = 1.7 MHz**

$$L_1 = 152 \text{ uH}$$

$$C_a = 200 \text{ pF (assumed)}$$

$$r = 30 \text{ ohms (assumed)}$$

$$A = 4.068 \times 10^5 \text{ ohms}$$

$$Q_1 Q_2 A = 1.29 \times 10^{10} \text{ ohms}$$

$$K = 8.324 \times 10^{-3}$$

$$\text{Check: } \frac{R_p}{1 + K^2 Q_1 Q_2} = 152 \text{ kohms} \gg \omega_r L_{eq} = 3.498 \text{ kohms}$$

### **Comments**

The values obtained for the coupling coefficient K hold for  $R_2 = \frac{R_{TNK}}{3}$ , as discussed previously. The transformed antenna-ground system resonance resistance, as seen from the secondary, will be equal to  $R_2$ , or  $\frac{R_{TNK}}{3}$ , as we are dealing with maximum power transfer to  $R_2$ . Under these conditions, the loaded Q of the secondary circuit will be 1/6 of the unloaded value. For other load conditions, the respective data should be entered into eq.(9).

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